An experimental investigation of interlaminar and intralaminar dynamic fracture of CFRPs: Effect of matrix modification using carbon nanotubes

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In this work, mode-I dynamic interlaminar and intralaminar fracture behaviors of carbon fiber reinforced polymers (CFRPs) are studied. Thick unidirectional composites were fabricated and their fracture performance was characterized under quasi-static three-point bending and dynamic one-point impact loading conditions. Both crack initiation and growth characteristics under stress-wave dominant conditions were evaluated in the latter case. The optical methods of digital image correlation (DIC) and ultra-high speed photography were employed to monitor crack tip deformations around transiently growing cracks. All relevant elastic properties were measured ultrasonically in order to determine stress intensity factors (SIFs). Interlaminar fracture responses were compared to the intralaminar counterparts using specimens of identical dimensions from the same original composite plate. Carbon nanotubes (CNTs) were then added with the aim of improving interlaminar fracture properties. While CNTs did not lead to improvements in critical stress intensity factor (\(K_{IC}\)), they did lead to modest improvements in fracture toughness (\(G_{IC}/G_{IC0}\)) under both quasi-static (+34%) and dynamic (+16%) loading conditions with significant scatter observed in these measurements.

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1. Introduction

In recent years carbon fiber reinforced polymers (CFRPs) have become a mainstay of aerospace structures. These layered systems are susceptible to fracture/damage in the interlaminar regions, particularly as a result of fatigue and/or impact loading; however, interlaminar fracture of composites is more tedious to characterize than intralaminar fracture. For a unidirectional composite, these two types of fracture are ideally very similar. Therefore, several authors [1–4] have compared intralaminar and interlaminar crack growth of unidirectional CFRPs under quasi-static loading conditions; however, to the authors’ knowledge, none have used the same specimen geometry when comparing intralaminar and interlaminar specimens. Most [1,2,4] used double cantilever beam (DCB) specimens to examine intralaminar fracture and compact tension specimens to examine interlaminar fracture, whereas de Moura et al. [3] used DCB specimens of different geometries to examine interlaminar and intralaminar fracture. Because measured stress intensity factors (SIFs) are dependent on specimen geometry [5–7], the present work involves the fabrication of thick carbon fiber composites such that both interlaminar and intralaminar specimens with the same geometry can be prepared from the same sheet and tested under similar conditions.

In order to reinforce the relatively weak interlaminar regions of CFRPs, several investigations [8–16] have added carbon nanotubes (CNTs) to this region. All used a DCB specimen to measure critical energy release rate (\(G_{IC}\)) under quasi-static loading conditions (typically following the ASTM Standard D5528 [8–10,12–16]). Some of the authors [9,10,13,15,16] formed their three-phase nanocomposites by dispersing CNTs into the resin first (Table 1), while others [8,11,12,14] began with carbon fiber sheets pre-impregnated with resin and then added CNTs by a sifting or spraying technique (Table 2). Most [8,11,12] of the latter group added CNTs only to the interlayer where the pre-crack would be introduced. Note that SWCNTs, DWCNTs, and MWCNTs refer to single-walled, double-walled, and multi-walled CNTs, respectively, whereas “UF” refers to unfunctionalized CNTs. In the current work, CNTs are incorporated first into the resin, and then the CNT/resin mixture is painted between layers of carbon fiber using a hand layup procedure similar to the methodology of Karapappas et al. [13] and Romhany and Szebenyi. [15]

Finally, the study of high loading-rate fracture in composites is critical for materials that will be used in aerospace applications.
where cold temperatures and high speeds typically elicit a more brittle response. Therefore, a few previous works [17–20] have employed strain gages in order to evaluate dynamic fracture parameters of interlaminar cracks in unidirectional CFRP samples. Instead of strain gages, the present work is the first to use digital image correlation (DIC) to map full-field deformations before and after crack initiation in order to determine interlaminar SIF histories in fiber-reinforced composite specimens subjected to dynamic impact loading. The advantage of using optical methods is that, unlike strain gages, they provide non-contact full-field deformations, as well as the precise location of the crack tip during the fracture event, which is necessary for estimating crack tip velocities for an accurate evaluation of dynamic SIF histories and hence crack growth resistance. This is particularly important because different CFRP systems could have different post-initiation resistance behaviors [21], an important factor in material selection.

In this context, it should be emphasized that energy release rate \( (G) \) is more relevant in characterizing fracture behavior of CFRPs. Further, direct evaluation of \( G \) without first finding SIFs is feasible under quasi-static conditions, as far-field load–displacement measurements are readily relatable to the crack tip stresses and displacements. However, the same is not true for stress-wave measurements: carbon fiber/epoxy (‘‘Neat’’) and CNT/carbon fiber/epoxy (with multi-walled NH\(_2\) functionalized NC 3152 CNTs from Nanocyl, <1 \( \mu \)m length, 9.5 nm diameter). CNTs were dispersed by calendering with a masterbatch technique using an Exakt 80E Calender. CNTs were first dispersed in the resin at 1.3 wt% using the calender until a thick paste was formed. This masterbatch was diluted with additional resin to 0.4 wt% CNTs by hand-stirring, followed by additional passes through the calender. Lastly, Epikure W was hand-stirred and mechanically stirred in, such that the final weight percentage of CNTs in epoxy was 0.3 wt%. Scanning electron microscopy of fracture surfaces of CNT/epoxy nanocomposites reported elsewhere [22] indicates good dispersion of CNTs by this method (Fig. 1a).

**2. Specimen preparation**

Unidirectional carbon fiber fabric was provided by V2 Composites, Inc. The resin system was Epon 862 and curing agent Epikure W from Momentive Specialty Chemicals, Inc. Two thick CFRPs were prepared in order to compare the effects of carbon nanotubes on the interlaminar and intralaminar fracture properties: carbon fiber/epoxy (‘‘Neat’’) and CNT/carbon fiber/epoxy (with multi-walled NH\(_2\)-functionalized NC 3152 CNTs from Nanocyl, <1 \( \mu \)m length, 9.5 nm diameter). CNTs were dispersed by calendering with a masterbatch technique using an Exakt 80E Calender. CNTs were first dispersed in the resin at 1.3 wt% using the calender until a thick paste was formed. This masterbatch was diluted with additional resin to 0.4 wt% CNTs by hand-stirring, followed by additional passes through the calender.

In light of this, all relevant in-plane and out-of-plane elastic properties necessary to evaluate both inter- and intralaminar SIFs are determined ultrasonically for each material system as part of the experimental program undertaken. This aspect is also unique to the reported dynamic fracture data as many other authors [18–21] adopted statically measured properties to estimate dynamic fracture parameters and/or relied on material characteristics of similar materials reported by others to accomplish the task. These authors [18–21] also used in-plane characteristics to estimate out-of-plane properties by assuming transverse isotropy.

In the present work, carbon nanotubes have been introduced into the interlaminar region of unidirectional CFRP specimens with the goal of enhancing interlaminar and intralaminar fracture properties under quasi-static and dynamic loading conditions. In order to achieve identical geometry between interlaminar and intralaminar fracture specimens, as well as to minimize the effects of wave reflections during dynamic fracture tests, sufficiently thick CFRP specimens have been fabricated. A methodology based on DIC has been developed to study dynamic interlaminar crack initiation and growth in a fiber-reinforced composite material subjected to dynamic loading. The SIF histories before \( a_0 \) and after crack initiation have been generated using displacement fields (determined by DIC) along with the ultrasonically-determined elastic properties.

### Table 2

<table>
<thead>
<tr>
<th>Author</th>
<th>(+%) G(_{\text{IC}})</th>
<th>Fiber Layout (manufacturing technique)</th>
<th>Fiber V(%)</th>
<th>CNT Type (dispersion technique)</th>
</tr>
</thead>
<tbody>
<tr>
<td>Almuhammadi et al. [8]</td>
<td>17</td>
<td>([0^\circ]_n) prepreg (CNTs sprayed on)</td>
<td>57</td>
<td>0.5 wt% COOH MWCNTs (sonication + solvent)</td>
</tr>
<tr>
<td>Hu et al. [11]</td>
<td>58</td>
<td>([0^\circ]_n) prepreg (CNTs silted onto interlayer)</td>
<td>57</td>
<td>10 g/m(^2) UF MWCNTs</td>
</tr>
<tr>
<td>Joshi and Dikshit [12]</td>
<td>40</td>
<td>Woven prepreg (CNTs sprayed on)</td>
<td>67</td>
<td>1.32 g/m(^2) UF SWCNTs (sonication + solvent)</td>
</tr>
<tr>
<td>Kim and Hahn [14]</td>
<td>6</td>
<td>([0^\circ]_n) prepreg (CNTs sprayed on)</td>
<td>65</td>
<td>0.5 wt% COOH SWCNTs (sonication + solvent)</td>
</tr>
</tbody>
</table>

\(a\) Indicates % improvement that was not significant relative to the reported error bars.

\(b\) Prepreg in this table indicates that CNTs were dispersed into the resin before the impregnation of dry fibers.
given in Fig. 4a. Three such interlaminar slices were stacked and glued together (Fig. 4b) in a vise using the same epoxy system as before in order to achieve ~50 mm thick interlaminar samples. Both interlaminar and intralaminar specimens were heated back to 120 °C for 2.5 h (in order to cure the epoxy glue used in interlaminar specimens) before post-curing for 3 h at 180 °C, such that the final geometry of all interlaminar and intralaminar fracture specimens were identical (Fig. 5), and all samples were subjected to identical curing schedules. Specimens were given a speckle coating for performing measurements using DIC. Pre-notches were then made with a diamond-tipped saw and sharpened by pressing a razor blade into the notch tip. This notching technique was chosen as an alternative to the difficulties of Teflon tape insertion for intralaminar samples detailed in Czabaj and Ratcliffe [2], who also measured interlaminar and intralaminar fracture of specimens machined from the same original composite sheet.

3. Experimental details

3.1. Quasi-static fracture test setup

Three-point bend quasi-static fracture tests (Fig. 6a) were performed on an Instron 4465 test stand at a crosshead speed of 0.01 mm/s using a 5 kN load cell. Load-point displacement was assumed to be the crosshead displacement. Load and displacement data were collected at a sample rate of 10 s⁻¹. Digital images were recorded every 3 s for determining SIFs from surface displacement fields. The critical SIF ($K_{IC}$) was chosen to be the SIF associated with the image immediately prior to visible crack tip movement. Crack growth was continued at least until the load decreased below half

![Fig. 1. Scanning electron micrograph of a fracture surface of an NH$_2$-CNT/epoxy nanocomposite showing good dispersion of CNTs.](image)

![Fig. 2. 0/90 scrim material (white) on the front and back surfaces of unidirectional carbon fiber fabric. Carbon fibers are running vertically.](image)

![Fig. 3. CFRP plate preparation details: (a) wet layup procedure for 60-layer CFRPs. (b) Orientation of interlaminar and intralaminar slices machined from the same thick composite plate.](image)

![Fig. 4. Interlaminar specimen preparation details: (a) Dimensions of a single interlaminar slice machined from the CFRP manufactured in Fig. 3 for dynamic fracture experiments. (b) Preparation of dynamic interlaminar fracture specimens by stacking and gluing three blocks with the dimensions shown in (a).](image)
of the peak load, at which time cracks had all grown well beyond an $a/W$ value of 0.5.

3.2. Dynamic fracture test setup

Dynamic 1-point impact tests were performed using a Hopkinson pressure bar (long-bar), depicted in Fig. 6b. Compressed air was used to launch a striker (25.4 mm diameter, 0.30 m length) coaxially towards the long-bar (25.4 mm diameter, 1.83 m length) of the same material (strain-rate independent aluminum 7075 T6). A soft aluminum pulse shaper was used between the striker and the long-bar to temper the loading rate such that more images could be captured in the pre- and post-initiation phases of the dynamic fracture event. The long-bar has a semi-circular (cylindrical) head machined on the end initially in contact with the specimen in order to deliver thickness-wise line-loading (1-point impact) to the specimen. The compressive stress wave traveled across the specimen, reached the free edge with the pre-notch, reflected back as a tensile wave, opened the pre-notch, initiated crack growth, and drove the crack to velocities as high as 800 m/s, all while inertia was holding the specimen within the view of the ultra-high speed camera. A roller of putty was placed above and below the specimen for the purpose of specimen alignment and in order to ensure symmetric wave reflections from the top and bottom surfaces of the specimen under “free-free” conditions. A Cordin ultra-high speed camera collected 32 digital images during the fracture event at 300,000 frames per second. A strain gage on the long-bar was used to collect stress wave data for independent estimation of SIFs prior to crack initiation using finite element analysis (FEA) software (see Section 3.5).

3.3. Experimental analysis

In both quasi-static and dynamic fracture experiments, undeformed images (recorded prior to loading) were paired with deformed images, and the resulting displacement fields were determined using ARAMIS image analysis software (DIC). Recorded 1000 by 1000 pixel images were divided into 15 by 15 pixel non-overlapping subimages (Fig. 7a), resulting in ~60 by 60 data points of displacement for each pair of images. All images have been flipped in order to comply with standard linear elastic fracture mechanics notation with origin at the crack tip and the crack growing in the positive x-direction. The location of the center of each subimage was tracked, and corresponding opening ($v$) and sliding ($u$) displacements were measured relative to their original positions in the undeformed images (Fig. 7b).

Sample digital images from an interlaminar dynamic fracture test of neat epoxy/carbon fiber are given in Fig. 8a, along with corresponding $v$-displacement fields (Fig. 8b) and $u$-displacement fields (Fig. 8c) determined by DIC. Notice that there are no discontinuities in the displacements where interlaminar specimens were glued together. Displacement data from the area behind the crack tip (where opening displacements are significant) in the range $0.5 < r/B < 1.5$, where $r$ is the radial distance of the polar coordinate system $(r, \theta)$ with origin at the crack tip, is used to determine SIFs in order to avoid three-dimensional (3-D) effects near the crack tip.

The following orthotropic displacement field equations [23] are used to determine SIFs for each deformed image for all quasi-static specimens, as well as for all dynamic specimens prior to crack initiation:

$$u(r, \theta) = K_I \frac{2r}{\pi} \left[ \frac{1}{2} \right] \left[ \begin{array}{c} \frac{1}{l_2 - l_1} \left( p_h l_2 z_1 - p_z l_2 z_2 \right) \\ \frac{1}{l_2 - l_1} \left( p_h l_2 z_1 - p_z l_2 z_2 \right) \end{array} \right]$$

$$v(r, \theta) = K_{II} \frac{2r}{\pi} \left[ \begin{array}{c} \frac{1}{l_2 - l_1} \left( q_h l_2 z_1 - q_z l_2 z_2 \right) \\ \frac{1}{l_2 - l_1} \left( q_h l_2 z_1 - q_z l_2 z_2 \right) \end{array} \right]$$

For intralaminar fracture of an orthotropic material with $x$–$z$ and $y$–$z$ planes as symmetry planes: $\mu_j (j = 1, 2)$ are the two roots of:

$$S_{11} \mu^4 + (2S_{12} + S_{66}) \mu^2 + S_{22} = 0,$$

where

$$p_h = \mu_1^1 S_{11} + S_{12},$$

$$q_h = S_{12} \cos \theta + \mu_1 \sin \theta,$$

$$\mu_1 = \sqrt{\cos \theta + \mu_1 \sin \theta},$$

$$\frac{1}{\mu_1} = \frac{1}{S_{11}} \frac{S_{22} - S_{12}}{S_{66}}$$

$$S_{11} S_{12} S_{13} 0 0 0$$

$$S_{12} S_{22} S_{23} 0 0 0$$

$$S_{13} S_{23} S_{66} 0 0 0$$

$$\frac{1}{\mu_2} = \frac{1}{S_{11}} \frac{S_{22} - S_{12}}{S_{66}}$$

$$S_{11} S_{12} S_{13} 0 0 0$$

$$S_{12} S_{22} S_{23} 0 0 0$$

$$S_{13} S_{23} S_{66} 0 0 0$$

$$\frac{1}{\mu_2} = \frac{1}{S_{11}} \frac{S_{22} - S_{12}}{S_{66}}$$

$$S_{11} S_{12} S_{13} 0 0 0$$

$$S_{12} S_{22} S_{23} 0 0 0$$

$$S_{13} S_{23} S_{66} 0 0 0$$

Equation (2) can be readily modified for the case of interlaminar fracture of an orthotropic material with $x$–$z$ and $y$–$z$ planes as symmetry planes, where $S_{13}$, $S_{23}$, and $S_{66}$ are used instead of $S_{13}$, $S_{23}$, and $S_{66}$, respectively.

In the above equations, $K_I$ is the mode-I SIF, $K_{II}$ is the mode-II SIF (which is expected to be near zero for this symmetric loading case), and $u$ and $v$ are the sliding and opening displacements. A least-squares analysis was used to determine a single $K_I$ and $K_{II}$ value for each pair of images. For consistency, the critical SIF ($K_c$ for quasi-static loading and $K_{c1D}$ for dynamic loading) was chosen to be the mode-I SIF value for the image immediately prior to visible crack initiation in the images. The $S$-matrix in Eq. (2) relates stresses and strains for an orthotropic material. Details of the

Fig. 5. Dimensions of fracture specimens subjected to (a) quasi-static loading and (b) dynamic impact loading. White dash lines indicate epoxy glue lines for interlaminar specimens. All specimens were given a speckle coating for performing DIC.
ultrasonic method for determining these elastic orthotropic material properties are given in Section 3.4. For the case of a dynamically growing crack, the orthotropic displacement field equations [24] are:

\[
u(r, \theta) = \frac{2}{C_{66}(c)} \text{Re} \left[ \left( \frac{\mu_2}{\lambda_1 - \lambda_2} \right) \left( \frac{\mu_1}{\lambda_1 - \lambda_2} \right) \sqrt{2\pi} \left( \frac{\mu_1 - \lambda_1}{\lambda_1 - \lambda_2} \right) \sqrt{\frac{z_2}{2\pi}} \right] K_1
+ \left( \frac{\gamma - \lambda \eta^2 \mu_2}{\lambda_1 - \lambda_2} \right) \sqrt{2\pi} \left( \frac{\mu_1 - \lambda_1}{\lambda_1 - \lambda_2} \right) \sqrt{\frac{z_2}{2\pi}} K_2,
\]

\[
u(r, \theta) = \frac{2}{C_{66}(c)} \text{Re} \left[ \left( \frac{\mu_2}{\lambda_1 - \lambda_2} \right) \left( \frac{\mu_1}{\lambda_1 - \lambda_2} \right) \sqrt{2\pi} \left( \frac{\mu_1 - \lambda_1}{\lambda_1 - \lambda_2} \right) \sqrt{\frac{z_2}{2\pi}} \right] K_1
+ \left( \frac{\gamma - \lambda \eta^2 \mu_2}{\lambda_1 - \lambda_2} \right) \sqrt{2\pi} \left( \frac{\mu_1 - \lambda_1}{\lambda_1 - \lambda_2} \right) \sqrt{\frac{z_2}{2\pi}} K_2.
\]

For intralaminar fracture of an orthotrop material with x–z and y–z planes as symmetry planes, \(\mu_j (j = 1, 2)\) are the 2 roots of:

\[
\mu^4 + \left( \frac{2\eta^2(c)}{c_1} \right) \mu^2 + \frac{2\eta^2(c)}{c_1^2} = 0,
\]

\[
\lambda_j = \frac{S_{11}}{S_{22}}, \quad \zeta = \frac{2S_{12} + S_{66}}{2\sqrt{S_{11}S_{22}}}, \quad \kappa = \frac{3\sqrt{S_{11}S_{22}} + S_{12}}{S_{11}S_{22} - S_{12}}, \quad c_1 = \sqrt{\frac{C_{11}}{\rho}},
\]

\[
c_4 = \frac{C_{66}}{\rho}, \quad \alpha_4(c) = 1 - \left( \frac{c_4(c)}{c_1} \right)^2, \quad \alpha_2^2(c) = 1 - \left( \frac{c_2(c)}{c_1} \right)^2,
\]

\[
\eta^2 = \left( \frac{\kappa + 1}{\kappa - 1} \right) \left( \frac{3 - \kappa + \zeta (\kappa + 1)}{4\sqrt{2}} \right), \quad \gamma = \sqrt{2} \eta^2 \left( \frac{3 - \kappa}{1 + \kappa} \right),
\]

\[
\lambda_j(c) = \frac{\eta^2 \alpha_4^2(c) + \eta^2 \alpha_2^2(c)}{1 + \eta^2 \alpha_2^2(c)} (j = 1, 2),
\]
are elements of the stiffness, 

\[ R(c) = \sqrt{\eta^2 q_i(c) q_j(c) - \frac{\eta q_j^2 q_i(c) + \eta^2 q_i^2(c)}{\sqrt{\eta^2 q_i(c) + q_i(c)}}} \]

\[ z_j = r(\cos \theta + \mu_j \sin \theta). \]

Again, for interlaminar fracture of an orthotropic material with \( x-z \) and \( y-z \) planes as symmetry planes, \( S_{11}, S_{22}, S_{55}, \) and \( S_{66} \) are used instead of \( S_{12}, S_{22}, S_{66}, \) and \( C_{66}, \) respectively, in Eq. (4).

In Eqs. (3) and (4), \( K_i \) and \( K_i' \) are the mode-I and mode-II SIFs for a dynamically growing crack tip, \( c \) is the crack tip velocity, \( \rho \) is the mass density, and \( C_{11} \) and \( C_{66} \) are elements of the stiffness (\( C \)) matrix, which is the inverse of the \( S \)-matrix given in Eq. (2):

\[
\begin{pmatrix}
\sigma_{11} \\
\sigma_{22} \\
\sigma_{33} \\
\sigma_{23} \\
\sigma_{31} \\
\sigma_{12}
\end{pmatrix} = \begin{pmatrix}
C_{11} & C_{12} & C_{13} & 0 & 0 & 0 \\
C_{21} & C_{22} & C_{23} & 0 & 0 & 0 \\
C_{31} & C_{32} & C_{33} & 0 & 0 & 0 \\
0 & 0 & 0 & C_{44} & 0 & 0 \\
0 & 0 & 0 & 0 & C_{55} & 0 \\
0 & 0 & 0 & 0 & 0 & C_{66}
\end{pmatrix} \cdot \begin{pmatrix}
\varepsilon_{11} \\
\varepsilon_{22} \\
\varepsilon_{33} \\
\varepsilon_{23} \\
\varepsilon_{31} \\
\varepsilon_{12}
\end{pmatrix}.
\]

\[ |C| = |S|^{-1}. \]

3.4. Ultrasonic determination of elastic constants

Because the primary emphasis of this work is dynamic fracture, ultrasonic determination of material properties is more appropriate than the quasi-static measurements typically reported in the literature. Several authors [25,26] have measured the elastic constants of composite materials; however, this will be the first work to measure the constants ultrasonically in both in-plane and out-of-plane directions for fracture parameter assessment.

The coefficients of the \( C \)-matrix were determined using an Epoch 600 Ultrasonic Flaw Detector from OLYMPUS in through-transmission mode. Composite material cubes were machined such that the faces aligned with the 1-, 2-, and 3-directions (Fig. 9), as defined originally in Fig. 3, where 1 – is the fiber and crack growth direction, 2 – is the crack opening direction for intralaminar specimens, and 3 – is the crack opening direction for interlaminar specimens.

Following the notations used in [27], let \( V_i \) denote the speed of a wave traveling in the \( i \)-direction with displacements in the \( j \)-direction. \( V_{11}, V_{22}, \) and \( V_{33} \) are then longitudinal wave speeds, whereas \( V_{12}, V_{13}, \) and \( V_{23} \) are shear wave speeds. These wave speeds are related to the diagonal terms of the \( C \)-matrix according to the following equations, where \( \rho \) is mass density:

\[
\begin{align*}
C_{11} &= \rho V_{11}^2, \quad C_{44} = \rho V_{44}^2 = \rho V_{22}^2, \\
C_{22} &= \rho V_{22}^2, \quad C_{55} = \rho V_{55}^2 = \rho V_{33}^2, \\
C_{33} &= \rho V_{33}^2, \quad C_{66} = \rho V_{66}^2 = \rho V_{11}^2.
\end{align*}
\]

In order to measure the off-diagonal terms of the \( C \)-matrix, three additional cubes were machined as before, and then parallel slices were taken with a 45° angle relative to the 1-, 2-, or 3-directions. A 45° rotation about the \( i \)-direction is shown in Fig. 10 with slices taken between the 2- and 3-faces. This new face is labeled the “6-face,” with corresponding longitudinal wave speed denoted as \( V_{66}. \)

Similarly, \( V_{66} \) was measured using opposing faces between the 1- and 2-faces, whereas \( V_{55} \) was measured using opposing faces.

\[ \text{For simplicity, the notation for wave speeds } V_{44}, V_{55}, \text{ and } V_{66} \text{ differs from [27].} \]
The off-diagonal C-matrix terms are then given by [27]:

\[
C_{12} = C_{21} = \sqrt{(C_{11} + C_{66} - 2\rho V_{44}^2)(C_{22} + C_{66} - 2\rho V_{44}^2) - C_{66}},
\]
\[
C_{13} = C_{31} = \sqrt{(C_{11} + C_{55} - 2\rho V_{55}^2)(C_{33} + C_{55} - 2\rho V_{55}^2) - C_{55}},
\]
\[
C_{23} = C_{32} = \sqrt{(C_{22} + C_{44} - 2\rho V_{66}^2)(C_{33} + C_{44} - 2\rho V_{66}^2) - C_{44}}.
\]

Once all 12 nonzero terms of the C-matrix coefficients are determined, the C-matrix (Eq. (5)) can be inverted in order to obtain the S-matrix (Eq. (2)). The resulting material properties are listed in Table 3, where “Neat” refers to epoxy/carbon fiber composites and “CNT” refers to CNT/epoxy/carbon fiber nanocomposites. The C-matrix coefficients are also compared to those given by Solodov et al. [26], who measured all 9 independent elastic constants ultra-
sonically for a nearly unidirectional carbon fiber composite (18 aligned 0° plies with 2 ± 45° plies in the center). Neat CFRP values follow the same trends as those found in the literature [26] at a slightly lower magnitude. This decrease in magnitudes may be due to a lower density of the material processed for this work. Additional application of this methodology to cortical bone is reported in Appendix A, where measured values performed by this methodology on bovine cortical bone closely matched those found in the literature (Table A1), further supporting the reported measurements.

Between the 1- and 3-faces. The off-diagonal C-matrix terms are then given by [27]:

\[
C_{12} = C_{21} = \sqrt{(C_{11} + C_{66} - 2\rho V_{44}^2)(C_{22} + C_{66} - 2\rho V_{44}^2) - C_{66}},
\]
\[
C_{13} = C_{31} = \sqrt{(C_{11} + C_{55} - 2\rho V_{55}^2)(C_{33} + C_{55} - 2\rho V_{55}^2) - C_{55}},
\]
\[
C_{23} = C_{32} = \sqrt{(C_{22} + C_{44} - 2\rho V_{66}^2)(C_{33} + C_{44} - 2\rho V_{66}^2) - C_{44}}.
\]

From the above table, a slight drop in the material density with the inclusion of CNTs into the matrix is evident. This is attributed to increased viscosity of CNT-modified resin, which caused less resin to be vacuumed out during curing. This contributed to modestly higher resin content, causing a decrease in \(E_1, E_2, G_{12}, G_{13},\) and fiber volume fraction \(V_f\). By intentionally manipulating the elastic constant input while solving Eqs. (1)-(4), it was found that the shear moduli (\(G_{12}, G_{13}\)) have the greatest effect on measured SIFs, whereas crack-opening direction elastic moduli (\(E_2, E_3\)) corresponding to intra- and inter-laminate crack opening directions, respectively, also make a significant contribution to the calculated SIF. Measured SIFs increase as shear moduli and elastic moduli in the respective crack opening directions increase. Meanwhile, the resulting fluctuations in Poisson’s ratios \(\nu_{12}, \nu_{13}\) have almost no effect on extracted SIFs. The reported Poisson’s ratios are less accurate than the elastic and shear moduli calculated by this ultrasonic method (Eqs. (6) and (7)). That is, a 1% change in ultrasonic wave speed \(V_{66}\) (easily within the uncertainty of this method) results in a 32% change in \(\nu_{12}\), a 15% change in \(\nu_{13}\), and a 1% change in \(E_1\). A similar 1% change in \(V_{11}\) results in a 2% change in \(E_1\), a 0.03% change in \(E_2\) and \(E_3\), and a 2% change in \(\nu_{12}\). Thus for applications of this method to materials of high degree of anisotropy where the accuracy of the Poisson terms is essential, the authors recommend independent verification of the Poisson’s ratios.

### Table 3

<table>
<thead>
<tr>
<th>Material property</th>
<th>Neat</th>
<th>CNT</th>
<th>C-matrix coefficient</th>
<th>Neat</th>
<th>Literature [26]</th>
</tr>
</thead>
<tbody>
<tr>
<td>(E_1) (GPa)</td>
<td>94.54</td>
<td>100.72</td>
<td>(C_{11}) (GPa)</td>
<td>102</td>
<td>127</td>
</tr>
<tr>
<td>(E_2) (GPa)</td>
<td>8.29</td>
<td>7.60</td>
<td>(C_{22}) (GPa)</td>
<td>11.6</td>
<td>13.8</td>
</tr>
<tr>
<td>(E_3) (GPa)</td>
<td>7.10</td>
<td>4.34</td>
<td>(C_{33}) (GPa)</td>
<td>10.0</td>
<td>12.8</td>
</tr>
<tr>
<td>(G_{12}) (GPa)</td>
<td>2.47</td>
<td>2.00</td>
<td>(C_{44}) (GPa)</td>
<td>2.5</td>
<td>3.6</td>
</tr>
<tr>
<td>(G_{13}) (GPa)</td>
<td>4.32</td>
<td>3.67</td>
<td>(C_{55}) (GPa)</td>
<td>4.3</td>
<td>5.0</td>
</tr>
<tr>
<td>(G_{23}) (GPa)</td>
<td>5.31</td>
<td>5.11</td>
<td>(C_{66}) (GPa)</td>
<td>5.3</td>
<td>6.7</td>
</tr>
<tr>
<td>(\nu_{12})</td>
<td>0.42</td>
<td>0.41</td>
<td>(C_{12}) (GPa)</td>
<td>8</td>
<td>7</td>
</tr>
<tr>
<td>(\nu_{13})</td>
<td>0.52</td>
<td>0.54</td>
<td>(C_{13}) (Gpa)</td>
<td>8</td>
<td>6</td>
</tr>
<tr>
<td>(\nu_{23})</td>
<td>0.54</td>
<td>0.58</td>
<td>(C_{23}) (GPa)</td>
<td>6</td>
<td>7</td>
</tr>
<tr>
<td>(\rho_{11})</td>
<td>0.04</td>
<td>0.03</td>
<td>(\rho ) (kg/m²)</td>
<td>1482</td>
<td>1600</td>
</tr>
<tr>
<td>(\rho_{12})</td>
<td>0.04</td>
<td>0.02</td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>(\rho_{13})</td>
<td>0.46</td>
<td>0.33</td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>Fiber (V_f)</td>
<td>50%</td>
<td>44%</td>
<td></td>
<td></td>
<td></td>
</tr>
</tbody>
</table>

\(\rho \) (kg/m²) | 1482 | 1442 |

**3.5 Finite element analysis**

Finite element analyses were carried out to supplement dynamic experimental results prior to crack initiation. A 3-D, transient, elasto-dynamic, finite element model using ABAQUS®/Explicit software was developed. The numerical model
included the specimen and the long-bar (Fig. 11) in order to ensure that the stress wave propagating into the specimen was captured as accurately as possible. Material property input for the model included all of the elastic properties in Table 3, along with the density. The model consisted of 218,000 tetrahedral elements with highly refined elements of size 0.1 mm in the impact and crack tip vicinities. The model had a total of 133,000 degrees of freedom. The particle velocity ($V_{pl}$) in the bar was determined from the measured strain history on the long-bar using Eq. (8)

$$V_{pl} = c \varepsilon$$

and input at the far right flat surface of the long-bar (Fig. 11). (In the above subscript $I$ denotes ‘incident,’ $c$ the bar wave speed and $\varepsilon$ the measured strain on the long-bar.) After propagating along the long-bar (most of which is not pictured in Fig. 11), the stress waves were transmitted into the specimen using a contact definition for the semi-circular region that impacts the flat edge of the specimen. Time increments automatically chosen during computations by the software were approximately 5 ns.

Instantaneous values of SIFs were computed from finite element results using a regression analysis of crack flank displacements. Apparent stress intensity factors ($K_I$) were calculated using crack opening displacements ($\delta_y$) in Eq. (9) near the crack tip according to:

$$\delta_y |_{y=\pm z, r=0} = 2R_{1} \sqrt{\frac{2 \tau}{\pi}} \Re \left[ \frac{1}{\mu_2 - \mu_1} (q_1 \mu_2 z_1 - q_2 \mu_1 z_2) \right].$$

The linear regions of $K_I$ versus $r$ plots were then extrapolated to the crack tip to obtain instantaneous SIFs as $K_I = \lim_{r \to 0} K$ [28,29].

4. Results

4.1. Quasi-static fracture results

To estimate baseline inter- and intra-laminar fracture parameters at crack initiation, quasi-static experiments were carried out on the 3-point bend configuration described in Section 3.1. Representative load–displacement curves are given for Neat (carbon fiber/epoxy without CNTs) samples in Fig. 12. Intralaminar specimens consistently showed more area under the curve than interlaminar specimens. Apart from the intrinsic inter- and intra-laminar elastic characteristics (Table 3), this increase is likely amplified by the presence of the thermoplastic-coated fiberglass scrim shown in Fig. 2 (initially believed by the authors to be an exclusively thermoplastic material). Thus, intralaminar cracks were aligned to propagate through the fiberglass scrim fibers, whereas interlaminar cracks propagated between layers of carbon fiber held by scrim without being affected by the scrim.

The measured SIF histories are shown in Fig. 13a and b from quasi-static fracture experiments for Neat samples (three experiments in each case to demonstrate repeatability) with cracks growing in the interlaminar and intralaminar directions, respectively. Intralaminar specimens had significantly higher critical SIFs than interlaminar specimens, likely attributed to the presence of the fiberglass scrim on each of the 60 layers of the original composite plates. Specimens exhibited controlled crack growth for several seconds after crack initiation (even longer for intralaminar specimens), and measured SIFs continued to increase due to fiber bridging [2]. The addition of CNTs (Fig. 13c and d) had no significant effect on quasi-static SIFs. The average quasi-static $K_{IC}$ values at crack initiation are summarized in Fig. 14, where CNTs led to no improvement in $K_{IC}$, whereas the intralaminar specimens had
significantly higher $K_{IC}$ values than the corresponding interlaminar specimens.

4.2. Dynamic fracture results

The measured SIF histories from dynamic fracture experiments are shown in Fig. 15a and b for Neat samples (three experiments in each case) with cracks growing in the interlaminar and intralaminar directions, respectively. As in the quasi-static case (Fig. 13), intralaminar specimens had significantly higher critical SIF ($K_d^{ill}$) and post-initiation SIFs ($K_d^i$) than interlaminar specimens, likely attributed to the presence of the thermoplastic-coated fiberglass scrim. The solid lines indicate FEA models (aligned with the average crack initiation SIF) generated using strain gage data on the long-bar which closely match the pre-initiation SIF values obtained from DIC and high speed photography. The addition of CNTs (Fig. 15c and d) gave similar $K_d^{ill}$ and $K_d^i$ values to corresponding Neat specimens, but with a higher degree of scatter. This increase in scatter was likely due to inconsistencies in dispersion of CNTs near the crack tip.

The measured crack tip velocity histories are given in Fig. 16. Following crack initiation, the intralaminar cracks quickly decelerated and eventually arrested, whereas the interlaminar cracks accelerated to 600–800 m/s. The intralaminar samples with CNTs had lower crack tip velocities than corresponding Neat specimens, indicating higher resistance to crack growth; however, these specimens also had reduced average $K_d^{ill}$ and $K_d^i$ values to corresponding Neat specimens, but with a higher degree of scatter.

The average dynamic $K_d^{ill}$ values are summarized in Fig. 17, where CNTs gave an insignificant reduction of $K_d^{ill}$, whereas the intralaminar specimens had significantly higher $K_d^{ill}$ values than the corresponding interlaminar specimens. All quasi-static (QS) and dynamic (D) critical SIF values
are listed in Table 4 along with the corresponding loading rate \( (K) \), quantified by the rate of change of SIF immediately before crack initiation.

Because fracture toughness is typically reported in terms of \( G_{IC} \) instead of critical SIF in the composites community, \( K_{IC} = \frac{K_d}{C_{0,ini}} \) values have been converted to \( G_{IC} = \frac{G_d}{C_{0,ini}} \) in Table 5 using the following equation for intralaminar fracture [30]:

\[
K_{IC} = \frac{K_d}{C_{0,ini}}
\]

\[
K_d = \frac{G_d}{C_{0,ini}}
\]
\[ G_I = -\frac{\mu}{2} \left( \frac{\mu_1 + \mu_2}{\mu_1} \right) \left( S_{11}\mu^4 + (2S_{12} + S_{66})\mu^2 + S_{22} \right) = 0 \]

and the following equation for interlaminar fracture:

\[ G_I = -\frac{\mu}{2} \left( \frac{\mu_1 + \mu_2}{\mu_1} \right) \left( S_{11}\mu^4 + (2S_{11} + S_{55})\mu^2 + S_{13} \right) = 0 \]

the average fracture toughness values for quasi-static \( G_{IC} \) and dynamic \( G_{ID} \) are given in Fig. 18, converted from critical SIFs according to Eqs. (10) and (11). While CNTs did not lead to improvements in critical stress intensity factors, they produced a 34% insignificant improvement in quasi-static interlaminar \( G_{IC} \) and a 16% insignificant improvement in dynamic interlaminar \( G_{ID} \) over Neat interlaminar samples. CNT-infused nanocomposites had similar interlaminar fracture toughness compared to corresponding neat samples, where fracture was dominated by the fiberglass scrim.

Table 4
Average quasi-static (QS) and dynamic (D) critical SIF values for interlaminar and intralaminar specimens.

<table>
<thead>
<tr>
<th>Test type</th>
<th>Average rate ( K ) (MPa ( \sqrt{\text{m}} ) s(^{-1} ))</th>
<th>Average Neat ( K_e/K_{IC}^{*} ) (MPa ( \sqrt{\text{m}} ))</th>
<th>Average CNT ( K_e/K_{IC}^{*} ) (MPa ( \sqrt{\text{m}} ))</th>
</tr>
</thead>
<tbody>
<tr>
<td>QS inter</td>
<td>1.31 ( \times 10^{-1} )</td>
<td>1.03 ( \pm 0.10 )</td>
<td>1.01 ( \pm 0.11 )</td>
</tr>
<tr>
<td>QS intra</td>
<td>1.76 ( \times 10^{-1} )</td>
<td>1.49 ( \pm 0.20 )</td>
<td>1.38 ( \pm 0.07 )</td>
</tr>
<tr>
<td>D inter</td>
<td>1.55 ( \times 10^{3} )</td>
<td>1.63 ( \pm 0.11 )</td>
<td>1.46 ( \pm 0.36 )</td>
</tr>
<tr>
<td>D intra</td>
<td>2.14 ( \times 10^{3} )</td>
<td>2.39 ( \pm 0.23 )</td>
<td>2.35 ( \pm 0.27 )</td>
</tr>
</tbody>
</table>

Table 5
Average fracture toughness \( G_{IC}/G_{ID}^{*} \) calculated from \( K_e/K_{IC}^{*} \) values (QS = quasi-static and D = dynamic).

<table>
<thead>
<tr>
<th>Test Type</th>
<th>Average Neat ( G_{IC}/G_{ID}^{*} ) (J/m(^2))</th>
<th>Average CNT ( G_{IC}/G_{ID}^{*} ) (J/m(^2))</th>
</tr>
</thead>
<tbody>
<tr>
<td>QS inter</td>
<td>109 ( \pm 22 )</td>
<td>149 ( \pm 31 )</td>
</tr>
<tr>
<td>QS intra</td>
<td>197 ( \pm 52 )</td>
<td>177 ( \pm 17 )</td>
</tr>
<tr>
<td>D inter</td>
<td>274 ( \pm 38 )</td>
<td>318 ( \pm 146 )</td>
</tr>
<tr>
<td>D intra</td>
<td>500 ( \pm 98 )</td>
<td>514 ( \pm 115 )</td>
</tr>
</tbody>
</table>

Fig. 18. Average fracture toughness values from critical SIFs using Eqs. (10) and (11) (a) quasi-static \( G_{IC} \) and (b) dynamic \( G_{ID} \) values. CNTs led to statistically insignificant improvements in both quasi-static (+34%) and dynamic (+16%) fracture toughness for interlaminar specimens, whereas interlaminar fracture of CNT-infused CFRPs, affected by fiberglass scrim, had similar fracture toughness compared to Neat samples.

Fig. 19. Scanning electron microscopy of (a) Neat CFRP at 1000\( \times \), (b) CNT CFRP at 4000\( \times \), and (c) CNT CFRP at 15000\( \times \). Crack growth is in the vertical direction. CNTs appear to be reasonably well-dispersed throughout the resin layers, but have little effect on resin layer fracture features or measured critical SIF.

Scanning electron micrographs of Neat and CNT samples are shown in Fig. 19. Crack growth occurred along the fiber direction, vertically in each image. Addition of well-dispersed CNTs as evident in Fig. 19b and c does not appear to affect fracture surface features in the resin layer compared to the Neat fracture surface in Fig. 19a. This observation is consistent with the lack of critical SIF enhancements in the reported measurements.
5. Discussion and conclusions

In this study, all interlaminar fracture specimens have higher initiation SIF under dynamic conditions than under quasi-static conditions, although this perceived loading rate effect could be due to the framing rate used as detailed in Appendix B. The quasi-static \( G_{IC} \) values for interlaminar fracture toughness in Table B1 of Appendix B are in good agreement with the literature [2,11,14,16].

Despite the high quality of dispersion demonstrated in Fig. 19, the addition of CNTs did not improve the critical stress intensity factor \( (K_{IC}/K_{IC,0}) \), although improvement (+34\%), comparable to those in the literature [8–16] in quasi-static interlaminar fracture toughness \( (G_{IC}) \) was found (Fig. 18). The errors are magnified in energy release rate calculations due to the squaring of the critical SIF in Eqs. (10) and (11), making improvements in fracture toughness statistically not significant. The scatter is particularly high for dynamic interlaminar fracture toughness of CNT-infused samples. Questions regarding scalability and cost-effectiveness of such preparations are not well documented at this time.

The volume fraction of fibers in the Neat CFRP processed for this work is on the lower end of the values reported in Tables 1 and 2 by other investigators. The addition of CNTs increased the viscosity of epoxy and caused a further reduction in \( V_f \), which led to a slight reduction in critical stress intensity factor under both quasi-static \( (K_{IC}) \) and dynamic \( (K_{IC,0}) \) loading conditions. In the ideal case, three-phase nanocomposites should be compared with Neat epoxy/carbon fiber composites of the same \( V_f \), although this variable is difficult to control due to the change in viscosity of the resin caused by well-dispersed CNTs.

The current work is unique in several respects:

1. The fabrication of relatively thick CFRPs allows for the comparison of interlaminar and intralaminar fracture data using the same testing procedures with samples machined to the same dimensions from the same original block of material. Under quasi-static loading conditions, where all material points sense the imposed loads simultaneously, measured SIFs are dependent on specimen geometry [5–7]. Under stress wave loading conditions, on the other hand, measured SIFs are also dependent on elastic wave speeds, which differ significantly between interlaminar and intralaminar directions (Table 3). The presence of fiberglass scrim on the front and back surfaces of the unidirectional carbon fiber fabric pinned intralaminar cracks whereas interlaminar cracks propagated through unsupported epoxy resin at the interlayer. Future work to achieve this idealized comparison of interlaminar and intralaminar fracture could involve similar preparation of thick carbon fiber composites using either unidirectional prepreg or dry unidirectional fabric held together by a thermoplastic scrim.

2. There is no current ASTM standard to study dynamic interlaminar fracture of composite materials. The current ASTM standard for quasi-static interlaminar fracture of composite materials (D5528) cannot be extended to study dynamic interlaminar fracture because thin specimens generally experience flexural stress waves. As the aerospace industry continues to replace aluminum with CFRP exposed to high speed events and cold temperatures, the behavior of these materials under dynamic loading conditions must be understood. Advantages of the current methodology involving DIC and high-speed photography include non-contact full-field deformations for the measurement of SIFs, as well as the precise location of the crack tip during the fracture event, which is necessary for estimating crack tip velocities for an accurate evaluation of dynamic SIF histories and hence crack growth resistance.

3. While several researchers [18–20] have previously investigated dynamic interlaminar fracture of CFRPs, they have assumed transverse isotropy and relied on statically-measured elastic properties in order to estimate dynamic fracture properties. In addition, they appear to have each used elastic properties reported by others who used similar materials. In other instances, out-of-plane elastic characteristics are estimated based on in-plane measurements of the same. The current research, on the other hand, includes the ultrasonic estimation of all dynamic elastic properties for orthotropic CFRP materials machined from the same composite plate, producing a high degree of consistency and reliability when in-plane and out-of-plane fracture characteristics are compared.

Acknowledgment

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Appendix A

In order to further verify the ultrasonic determination of elastic constants, the method was calibrated relative to data on bovine cortical bone (Fig. A1). Table A1 provides an overview of the studied elastic constants.

Table A1

<table>
<thead>
<tr>
<th>Material Property</th>
<th>Our Data (GPa)</th>
<th>Literature (GPa)</th>
</tr>
</thead>
<tbody>
<tr>
<td>( E_1 ) = ( E_L )</td>
<td>25.7</td>
<td>26.5</td>
</tr>
<tr>
<td>( E_2 ) = ( E_R )</td>
<td>20.6</td>
<td>18.1</td>
</tr>
<tr>
<td>( E_3 ) = ( E_T )</td>
<td>19.1</td>
<td>19.4</td>
</tr>
<tr>
<td>( G_{LT} )</td>
<td>8.17</td>
<td>7.22</td>
</tr>
<tr>
<td>( G_{LR} )</td>
<td>9.04</td>
<td>8.67</td>
</tr>
<tr>
<td>( G_{LT} )</td>
<td>9.49</td>
<td>8.65</td>
</tr>
<tr>
<td>( v_{12} ) = ( v_{RT} )</td>
<td>0.23</td>
<td>0.207</td>
</tr>
<tr>
<td>( v_{13} ) = ( v_{LT} )</td>
<td>0.32</td>
<td>0.305</td>
</tr>
<tr>
<td>( v_{23} )</td>
<td>0.31</td>
<td>0.325</td>
</tr>
<tr>
<td>( v_{13} ) = ( v_{LT} )</td>
<td>0.18</td>
<td>0.283</td>
</tr>
<tr>
<td>( v_{33} ) = ( v_{TT} )</td>
<td>0.24</td>
<td>0.285</td>
</tr>
<tr>
<td>( v_{12} )</td>
<td>0.29</td>
<td>0.222</td>
</tr>
</tbody>
</table>

\( \rho \) (kg/m\(^3\))
| 2000 | ~ |

Fig. A1. Alignment of cortical bone ultrasonic specimens.
cortical bone available in the literature. The bovine cortical bone was dried and machined to cubes of ~6 mm side, with each face aligned with one of the three material directions (see Fig. 10). Ultrasonic results were successfully compared to those reported by Van Buskirk et al. [27] from a dried human femur (Table A1), further confirming the method used for fiber reinforced composites work.

**Appendix B**

Using the current methodology, critical SIF refers to the SIF value calculated from the last image before the crack was observed to move in the sequence of photographs. Alternatively, if the initiation SIF is chosen to be the first image when the crack is observed to move, the assigned quasi-static initiation SIF values increase significantly, whereas the assigned dynamic initiation SIF values increase only slightly. This is because quasi-static SIFs continue to increase at approximately the same rate before and after crack initiation (Fig. 13), whereas dynamic SIFs level off dramatically after initiation (Fig. 15), particularly in the interlaminar case where the scrim fibers do not participate in crack growth. Accordingly, Table B1 compares Neat $KIC/KIC_{ini}$ and $GIC/GIC_{ini}$ values calculated from the image before and after crack initiation.

In light of this, future works could use a slower loading rate or a faster framing rate, particularly under quasi-static conditions, in order to improve estimation of $KIC$ and $GIC$ values by this methodology.

**References**


**Table B1**

<table>
<thead>
<tr>
<th>Test type</th>
<th>Neat before $KIC/KIC_{ini}$ (MPa √m)</th>
<th>Neat After $KIC/KIC_{ini}$ (MPa √m)</th>
<th>Neat before $GIC/GIC_{ini}$ (J/m²)</th>
<th>Neat after $GIC/GIC_{ini}$ (J/m²)</th>
</tr>
</thead>
<tbody>
<tr>
<td>QS inter</td>
<td>1.03 ± 0.10</td>
<td>1.57 ± 0.24</td>
<td>109 ± 22</td>
<td>256 ± 75</td>
</tr>
<tr>
<td>QS intra</td>
<td>1.49 ± 0.20</td>
<td>1.90 ± 0.19</td>
<td>197 ± 52</td>
<td>316 ± 65</td>
</tr>
<tr>
<td>D inter</td>
<td>1.63 ± 0.11</td>
<td>1.67 ± 0.09</td>
<td>274 ± 38</td>
<td>285 ± 32</td>
</tr>
<tr>
<td>D intra</td>
<td>2.39 ± 0.23</td>
<td>2.49 ± 0.22</td>
<td>500 ± 98</td>
<td>544 ± 95</td>
</tr>
</tbody>
</table>